### From the quantum Lie algebra $U_q(\mathfrak{sl}_2)$ to the ASEP(q, j)

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#### Outline

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- Constructive approach to duality theory via Lie algebra
- $\mathfrak{su}_q(2)$  algebra: construction of ASEP(q, j)
- Properties of ASEP(q, j)

1. Constructive approach

### to duality theory

via Lie algebra

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#### **Stochastic Duality**

 $(\eta_t)_{t\geq 0}$  Markov process on  $\Omega$  with generator *L*,

 $(\xi_t)_{t\geq 0}$  Markov process on  $\Omega_{dual}$  with generator  $L_{dual}$ 

 $\xi_t$  is dual to  $\eta_t$  with duality function  $D : \Omega \times \Omega_{dual} \to \mathbb{R}$  if  $\forall t \ge 0$  $\mathbb{E}_{\eta}(D(\eta_t, \xi)) = \mathbb{E}_{\xi}(D(\eta, \xi_t)) \qquad \forall (\eta, \xi) \in \Omega \times \Omega_{dual}$ 

 $\eta_t$  is self-dual if  $L_{dual} = L$ .

Duality is equivalent to  $LD(\cdot,\xi)(\eta) = L_{dual}D(\eta,\cdot)(\xi)$ 

Self-duality: (L = L<sub>dual</sub>) for a Markov chain with countable state space it is equivalent to

 $LD = DL^T$ 



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Reversibility and trivial self-duality: if μ is a reversible measure, a trivial (i.e. diagonal) self-duality function is

$${\sf d}(\eta,\xi)={1\over \mu(\eta)}\delta_{\eta,\xi}$$

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- Symmetries and (non-trivial) self-duality:
  - S: symmetry of the Markov generator, i.e.  $[\mathbf{L}, \mathbf{S}] = 0$
  - d: trivial self-duality function
  - $\longrightarrow$  **D** = **Sd** is a self-duality function

#### Lie algebra

A Lie algebra is a vector space  $\mathfrak{g}$  over a field F with a binary operation  $[\cdot, \cdot] : \mathfrak{g} \times \mathfrak{g} \to \mathfrak{g}$  (Lie bracket)

- $[\cdot, \cdot]$  is bilinear
- $\forall u, v \text{ in } \mathfrak{g}: [u, v] = -[v, u]$
- ▶ [Jacobi identity]:  $\forall u, v, w$  in  $\mathfrak{g}$

[u, [v, w]] + [w, [u, v]] + [v, [w, u]] = 0

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#### Algebraic approach

- 1. Write the Markov generator in abstract form, i.e. as an element of a Lie algebra, using the algebra generators.
- 2. Duality is related to a change of representation. Duality functions are the intertwiners.
- 3. Self-duality is associated to symmetries.

Conversely, the approach can be turned into a constructive method.

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vi) (*Markov generator*): Apply a "ground state transformation" to turn *H* into a Markov generator *L*.

## 2. $\mathfrak{su}_q(2)$ algebra:

## construction of ASEP(q, j)

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#### q-numbers

For  $q \in (0, 1)$  and  $n \in \mathbb{N}_0$  introduce the *q*-number

$$[n]_q = \frac{q^n - q^{-n}}{q - q^{-1}}$$

Remark:  $\lim_{q\to 1} [n]_q = n$ .

#### The first *q*-number's are:

 $[0]_q = 0,$   $[1]_q = 1,$   $[2]_q = q + q^{-1},$   $[3]_q = q^2 + 1 + q^{-2},$  ...

The quantum Lie algebra  $\mathfrak{su}_q(2) \equiv U_q(\mathfrak{sl}_2)$ 

For  $q \in (0, 1)$  consider the algebra with generators  $J^+, J^-, J^0$ 

 $[J^+, J^-] = [2J^0]_q, \qquad [J^0, J^{\pm}] = \pm J^{\pm}$ 

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The Casimir element:

$$C = J^{-}J^{+} + [J^{0}]_{q}[J^{0} + 1]_{q}$$

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A standard representation ( $n = 0, 1, \dots 2j$ )

$$\begin{cases} J^{+} |n\rangle &= \sqrt{[2j-n]_{q}[n+1]_{q}} |n+1\rangle \\ J^{-} |n\rangle &= \sqrt{[n]_{q}[2j-n+1]_{q}} |n-1\rangle \\ J^{0} |n\rangle &= (n-j) |n\rangle \end{cases}$$

In this representation  $C |n\rangle = [j]_q [j+1]_q |n\rangle$ 

#### Co-product

A coproduct on an algebra  $\mathfrak{g} \to \mathfrak{g} \otimes \mathfrak{g}$  is an algebra homomorphism:

 $\Delta([A,B]) = [\Delta(A), \Delta(B)] \quad \forall A, B \in \mathfrak{g}$ 



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Define the co-product  $\Delta : U_q(\mathfrak{sl}_2) \to U_q(\mathfrak{sl}_2)^{\otimes 2}$  as follows

$$egin{array}{rcl} \Delta(J^{\pm}) &=& J^{\pm}\otimes q^{-J^0}+q^{J^0}\otimes J^{\pm}\ \Delta(J^0) &=& J^0\otimes 1+1\otimes J^0 \end{array}$$

then

 $[\Delta(J^+),\Delta(J^-)] = [2\Delta(J^0)]_q \qquad [\Delta(J^0),\Delta(J^\pm)] = \pm \Delta(J^\pm)$ 

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#### Quantum Hamiltonian

$$\Delta(C_i) = -q^{J_i^0} \Big\{ J_i^+ \otimes J_{i+1}^- + J_i^- \otimes J_{i+1}^+ + B_{i,i+1} \Big\} q^{-J_{i+1}^0}$$

$$\begin{array}{lll} B_{i,i+1} & = & \displaystyle \frac{(q^{i}+q^{-i})(q^{i+1}+q^{-(j+1)})}{2(q-q^{-1})^{2}} \left(q^{J_{i}^{0}}-q^{-J_{i}^{0}}\right) \otimes \left(q^{J_{i+1}^{0}}-q^{-J_{i+1}^{0}}\right) \\ & + & \displaystyle \frac{(q^{i}-q^{-i})(q^{i+1}-q^{-(j+1)})}{2(q-q^{-1})^{2}} \left(q^{J_{i}^{0}}+q^{-J_{i}^{0}}\right) \otimes \left(q^{J_{i+1}^{0}}+q^{-J_{i+1}^{0}}\right) \end{array}$$

$$H_{(L)} := \sum_{i=1}^{L-1} \left( 1^{\otimes (i-1)} \otimes \Delta(C_i) \otimes 1^{\otimes (L-i-1)} + c_{q,j} 1^{\otimes L} \right)$$
$$(q^{2j} - q^{-2j})(q^{2j+1} - q^{-(2j+1)}) \qquad \text{a.t.} \quad H(\mathbb{Q}^{L} \mid 0))$$

$$c_{q,j} = \frac{(q^{L} - q^{-1})(q^{L} - q^{-1})}{(q - q^{-1})^{2}} \qquad s.t. \quad H\left(\otimes_{i=1}^{L} |0\rangle\right) = 0$$

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#### Symmetries of H

Iterate the co-product  $\Delta^n: U_q(\mathfrak{sl}_2) \to U_q(\mathfrak{sl}_2)^{\otimes (n+1)}$ , i.e. for  $n \geq 2$ 

$$\begin{array}{lll} \Delta^n(J^{\pm}) &=& \Delta^{n-1}(J^{\pm}) \otimes q^{-J^0} + q^{\Delta^{n-1}(J^0)} \otimes J^{\pm} \\ \Delta^n(J^0) &=& \Delta^{n-1}(J^0) \otimes 1 + 1^{\otimes n} \otimes J^0 \end{array}$$

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#### Lemma

$$\mathcal{J}^{\pm} := \Delta^{L-1}(J^{\pm}) = \sum_{i=1}^{L} q^{J_{1}^{0}} \otimes \cdots \otimes q^{J_{i-1}^{0}} \otimes J_{i}^{\pm} \otimes q^{-J_{i+1}^{0}} \otimes \ldots \otimes q^{-J_{L}^{0}}$$
$$\mathcal{J}^{0} := \Delta^{L-1}(J^{0}) = \sum_{i=1}^{L} \underbrace{1 \otimes \cdots \otimes 1}_{(i-1) \text{ times}} \otimes J_{i}^{0} \otimes \underbrace{1 \otimes \cdots \otimes 1}_{(L-i) \text{ times}}.$$

are symmetries of  $H_{(L)}$ .

not a stochastic generator!

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• construct a non-trivial symmetry of *H* from the trivial ones:  $\mathcal{J}^{\pm,0}$ 

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- ► construct a non-trivial symmetry of *H* from the trivial ones: J<sup>±,0</sup>
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► construct a non-trivial symmetry of L using the fact that if S is a symmetry of H then G<sup>-1</sup>SG is a symmetry of L

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Strategy

- ► construct a non-trivial symmetry of *H* from the trivial ones: J<sup>±,0</sup>
- transform H into a stochastic generator L via a transformation

 $L = G^{-1}HG$ 

- ► construct a non-trivial symmetry of L using the fact that if S is a symmetry of H then G<sup>-1</sup>SG is a symmetry of L
- use the non-trivial symmetries of L to construct self duality functions for the associated Markov process

#### Ground State transformation

#### Lemma

Let *H* be a matrix with  $H(\eta, \eta') \ge 0$  for  $\eta \ne \eta'$ . Suppose *g* is a positive ground state. i.e. Hg = 0 and  $g(\eta) > 0$ . Let *G* be the matrix  $G(\eta, \eta') = g(\eta)\delta(\eta, \eta')$ . Then

 $L = G^{-1} H G$ 

is a Markov generator.

Proof.

$$L(\eta,\eta') = rac{H(\eta,\eta')g(\eta')}{g(\eta)}$$

Therefore

$$L(\eta,\eta') \ge 0$$
 if  $\eta \ne \eta'$ 

$$\sum_{\eta'} L(\eta, \eta') = \mathbf{0}$$

#### Exponential symmetries

- $g^{(0)} = \bigotimes_{i=1}^{L} |0\rangle$  is a ground state, i.e.  $Hg^{(0)} = 0$ .
- For every symmetry [H, S] = 0 another ground state is  $g = Sg^{(0)}$ .
- The exponential symmetry

$$S^+ = \exp_{q^2}(E) = \sum_{n \ge 0} \frac{(E)^n}{[n]_q!} q^{-n(n-1)/2}$$

where

$$E = \Delta^{(L-1)}(q^{J^0} J^+)$$

gives a positive ground state

$$g = \mathcal{S}^+ g^{(0)} = \sum_{\ell_1,...,\ell_L} \otimes_{i=1}^L \left( \sqrt{\binom{2j}{\ell_i}_q} \cdot q^{\ell_i(1+j-2ji)} 
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#### ASEP(q,j) process

#### Definition

The Markov process ASEP(q, j) on  $[1, L] \cap \mathbb{Z}$ , denoted by  $(\eta(t))_{t \ge 0}$ , with state space  $\{0, 1, \dots, 2j\}^{L}$  is defined by

$$(L^{ASEP(q,j)}f)(\eta) = \sum_{i=1}^{L-1} (L_{i,i+1}f)(\eta)$$

#### with

$$\begin{aligned} (L_{i,i+1}f)(\eta) &= q^{\eta_i - \eta_{i+1} - (2j+1)} [\eta_i]_q [2j - \eta_{i+1}]_q (f(\eta^{i,i+1}) - f(\eta)) \\ &+ q^{\eta_i - \eta_{i+1} + (2j+1)} [2j - \eta_i]_q [\eta_{i+1}]_q (f(\eta^{i+1,i}) - f(\eta)) \end{aligned}$$

Remark: it follows from  $L = G^{-1}HG$ .

#### ASEP(q,j) process: special cases

$$\begin{aligned} (L_{i,i+1}f)(\eta) &= q^{\eta_i - \eta_{i+1} - (2j+1)} [\eta_i]_q [2j - \eta_{i+1}]_q (f(\eta^{i,i+1}) - f(\eta)) \\ &+ q^{\eta_i - \eta_{i+1} + (2j+1)} [2j - \eta_i]_q [\eta_{i+1}]_q (f(\eta^{i+1,i}) - f(\eta)) \end{aligned}$$

- q = 1 → SEP(j): symmetric partial exclusion jump right at rate η<sub>i</sub>(2j − η<sub>i+1</sub>), jump left at rate (2j − η<sub>i</sub>)η<sub>i+1</sub>
- ►  $j = 1/2 \rightarrow ASEP(q)$ : asymmetric exclusion jump right at rate  $q^{-1}$ , jump left at rate q
- ►  $j = \infty \rightarrow TAZRP(q)$ : totally asymmetric zero range after rescaling time  $t \rightarrow q^{4j-1}t$ , jump right at rate  $\frac{1-q^{2\eta_j}}{1-q^2}$

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## 3. Properties of ASEP(q, j)

#### Properties of ASEP(q,j)

Theorem

a) The ASEP(q, j) is well-defined on  $\mathbb{Z}$  and is a monotone process.

b) The ASEP(q, j) on  $\mathbb{Z}$  has a family (labeled by  $\alpha > 0$ ) of reversible product measures with marginals

$$\mathbb{P}_{\alpha}(\eta_{i} = x) = \frac{\alpha^{x}}{Z_{i,\alpha}} {\binom{2j}{x}}_{q} \cdot q^{2x(1+j-2ji)}$$

c) The ASEP(q, j) has translation invariant stationary product measures only for j = 1/2 and for  $j \to \infty$ .

Proof of b) for  $\alpha = 1$ : using  $H = H^T$ 

 $G^{2}L = (G^{2}G^{-1}HG) = (GHG^{-1})G^{2} = L^{T}G^{2}$ 

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#### Self-duality of ASEP(q, j)

Theorem The ASEP(q, j) on  $\mathbb{Z}$  is self-dual on

$$D(\eta,\xi) = \prod_{i=1}^{L} \frac{[\eta_i]_q!}{[\eta_i - \xi_i]_q!} \frac{\Gamma_q(2j+1-\xi_i)}{\Gamma_q(2j+1)} \cdot q^{(\eta_i - \xi_i)[2\sum_{k=1}^{i-1} \xi_k + \xi_i] + 4ji\xi_i} \cdot 1_{\eta_i \ge \xi_i}$$

Proof: it follows from the general method

- $d = G^{-2}$  is a trivial duality function
- $[H, S^+]$  and  $L = G^{-1}HG$ , thus  $[L, G^{-1}S^+G] = 0$ .
- $D = (G^{-1}S^+G)G^{-2} = G^{-1}S^+G^{-1}$  is a duality fct.

#### Current of ASEP(q, j)

### Definition

Let

$$N_i(t) := \sum_{k \ge i} \eta_k(t)$$

The current  $J_i(t)$  during the time interval [0, t] across the bond (i - 1, i) is defined as the net number of particles traversing the bond in the right direction:

$$J_i(t)=N_i(t)-N_i(0)$$

Remark: let  $\xi^{(i)}$  be the configuration with 1 dual particle:

$$\xi_m^{(i)} = \begin{cases} 1 & \text{if } m = i \\ 0 & \text{otherwise} \end{cases}$$

then

$$D(\eta,\xi^{(i)}) = \frac{q^{4ji-1}}{q^{2j}-q^{-2j}} \cdot (q^{2N_i}-q^{2N_{i+1}})$$

#### First $q^2$ -moment of the current

Theorem

$$\mathbb{E}_{\eta} \left[ q^{2J_{i}(t)} \right] = q^{\sum_{k < i} \eta_{k}} \\ - \sum_{k < i} q^{-4jk} \mathbb{E} \left[ q^{4jX(t)} \left( 1 - q^{-2\eta_{X(t)}} \right) q^{2(N_{X(t)}(0) - N_{i}(0))} \mid X(0) = k \right]$$

with X(t) a random walker on  $\mathbb{Z}$  jumping left at rate  $q^{2j}[2j]_q$  and jumping right at rate  $q^{-2j}[2j]_q$ 

$$\mathbb{P}(X(t) = z \mid X(0) = k) = e^{-[4j]_q t} q^{-2j(z-k)} I_{z-i}(2[2j]_q t)$$

 $I_n(t)$  modified Bessel fct.

#### First $q^2$ -moment of the current

Proof: Duaility gives

 $\mathbb{E}_{\eta}(D(\eta(t),\xi^{(i)})) = \mathbb{E}_{\xi^{(i)}}(D(\eta,\xi^{(X(t))})$ 

$$\mathbb{E}_{\eta}\left[q^{4ji} \cdot (q^{2N_{i}(t)} - q^{2N_{i+1}(t)})\right] = \mathbb{E}_{\xi^{(i)}}\left[q^{4jX(t)} \cdot (q^{2N_{X(t)}} - q^{2N_{X(t)+1}})\right]$$

Therefore

$$\mathbb{E}_{\eta}[q^{2N_{i}(t)}] = \mathbb{E}_{\eta}[q^{2N_{i+1}(t)}] + q^{-4ji}\mathbb{E}_{\xi^{(i)}}\left[q^{4jX(t)} \cdot (q^{2N_{X(t)}} - q^{2N_{X(t)+1}})\right]$$

Multiply by  $q^{-2N_i(0)}$  to get a recursion relation for the current and iterate.

#### Step initial condition

Proposition

For the step initial conditions  $\eta^\pm \in \{\mathbf{0},\ldots,\mathbf{2j}\}^\mathbb{Z}$  defined as

$$\eta_i^+ := \left\{ egin{array}{cccc} 0 & ext{for} & i < 0 \ 2j & ext{for} & i \geq 0 \end{array} 
ight. \qquad \eta_i^- := \left\{ egin{array}{ccccc} 2j & ext{for} & i < 0 \ 0 & ext{for} & i \geq 0 \end{array} 
ight.$$

one has

$$\mathbb{E}_{\eta^{+}}\left[q^{2J_{i}(t)}\right] = q^{4j\max\{0,i\}} \left\{1 + q^{-4ji} \mathbf{E}_{i}\left[\left(1 - q^{4jX(t)}\right) \mathbf{1}_{X(t)\geq 1}\right]\right\}$$
$$\mathbb{E}_{\eta^{-}}\left[q^{2J_{i}(t)}\right] = q^{-4j\max\{0,i\}} \left\{1 - \mathbf{E}_{i}\left[\left(1 - q^{4jX(t)}\right) \mathbf{1}_{X(t)\geq 1}\right]\right\}$$

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#### Step initial condition

#### Remark 1: asymptotics

$$\begin{split} &\lim_{t\to\infty} \mathbb{E}_{\eta^+} \left[ q^{2J_i(t)} \right] = q^{4j \max\{0,i\}} \left( 1 + q^{-4ji} \right) \qquad \text{shock} \\ &\lim_{t\to\infty} \mathbb{E}_{\eta^-} \left[ q^{2J_i(t)} \right] = 0 \qquad \text{rarefaction fan} \end{split}$$

#### Remark 2: contour integral

$$\mathbb{E}_{\eta^{+}}\left[q^{2J_{k}(t)}\right] = \frac{q^{4j\max\{0,k\}}}{2\pi i} \int e^{-\frac{q^{2j}[2j]_{q}^{3}(q^{-1}-q)^{2}z}{(1+q^{4j}z)(1+z)}t} \left(\frac{1+z}{1+q^{4j}z}\right)^{k} \frac{dz}{z}$$

where the integration contour includes 0 and  $-q^{-4j}$  but does not include -1.